

**Hessenberg eigenvalue – eigenvector relations  
and their application to the error analysis of  
finite precision Krylov subspace methods**

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Minisymposium on Numerical Linear Algebra  
Technical University Hamburg–Harburg  
10.07.2002

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# The Menagerie of Krylov Methods

- o Lanczos based methods (short-term methods)
- o Arnoldi based methods (long-term methods)
  
- o eigensolvers:  $Av = v\lambda$
- o linear system solvers:  $Ax = b$ 
  - o (quasi-) orthogonal residual approaches: (Q)OR
  - o (quasi-) minimal residual approaches: (Q)MR

Extensions:

- o Lanczos based methods:
  - o look-ahead
  - o product-type (LTPMs)
  - o applied to normal equations (CGN)
- o Arnoldi based methods:
  - o restart (thin/thick, explicit/implicit)
  - o truncation (standard/optimal)

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## A Unified Matrix Description of Krylov Methods

Krylov methods as **projection** onto 'simpler' matrices:

- $Q^H Q = I$ ,  $Q^H A Q = H$  Hessenberg (Arnoldi),
- $\hat{Q}^H Q = I$ ,  $\hat{Q}^H A Q = T$  tridiagonal (Lanczos)

Introduce computed (condensed) matrix  $C = T, H$

$$Q^{-1} A Q = C \quad \Rightarrow \quad A Q = Q C$$

Iteration implied by **unreduced Hessenberg** structure:

$$A Q_k = Q_{k+1} \underline{C}_k, \quad Q_k = [q_1, \dots, q_k], \quad \underline{C}_k \in \mathbb{K}^{(k+1) \times k}$$

Stewart: '**Krylov Decomposition**'

Iteration spans Krylov subspace ( $q = q_1$ ):

$$\text{span}\{Q_k\} = \mathcal{K}_k = \text{span}\{q, Aq, \dots, A^{k-1}q\}$$

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## Perturbed Krylov Decompositions

A Krylov decomposition **analogue** holds true in **finite precision**:

$$\begin{aligned}AQ_k = Q_{k+1}\underline{C}_k - F_k &= Q_k C_k + q_{k+1} c_{k+1,k} e_k^T - F_k \\ &= Q_k C_k + M_k - F_k\end{aligned}$$

We have to investigate the impacts of the method on

- o the **structure** of the **basis**  $Q_k$  (local orthogonality/duality)
- o the **structure** of the **computed**  $C_k, \underline{C}_k$
- o the **size/structure** of the **error term**  $-F_k$

**Convergence** theory:

- o is usually based on **inductively** proven properties:  
orthogonality, bi-orthogonality,  $A$ -conjugacy, ...

What can be said about these properties?

'Standard' **error analysis**:

- o splits into *forward* and *backward* error analysis.

Does this analysis apply to Krylov methods?

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All methods fit **pictorially** into:

The diagram shows a matrix equation represented by boxes and symbols. On the left, a large square box labeled  $A$  is followed by a vertical rectangular box labeled  $Q_k$ . A minus sign is placed between them. To the right of the minus sign is another vertical rectangular box labeled  $Q_k$ . Above this second  $Q_k$  box is a small square box labeled  $C_k$ . An equals sign follows. To the right of the equals sign is a vertical rectangular box labeled  $0$ . A minus sign is placed between the  $0$  box and a final vertical rectangular box labeled  $F_k$ . A period follows the  $F_k$  box.

This is a perturbed Krylov decomposition, as **subspace equation**.

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- Examination** of the methods has to be done according to
- o methods **directly based on the Krylov decomposition**
  - o methods **based on a split Krylov decomposition**
  - o **LTPMs**

The **matrix  $C_k$**  plays a crucial role:

- o  $C_k$  is **Hessenberg** or even **tridiagonal** (**basics**),
- o  $C_k$  may be **blocked** or **banded** (**block Krylov methods**),
- o  $C_k$  may have **humps, spikes, ...** (**more sophisticated**)

The **error analysis** and **convergence theory** splits further up:

- o knowledge on **Hessenberg (tridiagonal)** matrices
- o knowledge on **orthogonality, duality, conjugacy, ...**

We start with results on Hessenberg matrices.

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## A Short Excursion on Matrix Structure

$J_\Lambda$	Jordan matrix of $A$
$V$	right eigenvector-matrix, $AV = VJ_\Lambda$
$\hat{V} \equiv V^{-H}$	left eigenvector-matrix, $\hat{V}^H A = J_\Lambda \hat{V}^H$
$\check{V} \equiv V^{-T}$	alternate eigenvector-matrix, $\check{V}^T A = J_\Lambda \check{V}^T$
$\chi_A(\lambda) \equiv \det(\lambda I - A)$	characteristic polynomial of $A$
$R(\lambda) \equiv (\lambda I - A)^{-1}$	resolvent
$C_k(A)$	$k$ th compound matrix of $A$
$\text{adj}(A)$	classical adjoint, adjugate of $A$
$A_{ij}$	$A$ with row $i$ and column $j$ deleted
$S, S_i$	sign matrices

The **adjoint** of  $\lambda I - A$  fulfils

$$\text{adj}(\lambda I - A)(\lambda I - A) = \det(\lambda I - A)I.$$

Suppose that  $\lambda$  is not contained in the spectrum of  $A$ .



We form the resolvent of  $\lambda$  and obtain

$$\begin{aligned} \text{adj}(\lambda I - A) &= \det(\lambda I - A) R(\lambda) \\ &= V \left( \chi_A(\lambda) J_{\lambda-\Lambda}^{-1} \right) \widehat{V}^H. \end{aligned}$$

The shifted and inverted Jordan matrix looks like

$$J_{\lambda-\lambda_i}^{-1} = S_i E_i S_i \equiv S_i \begin{pmatrix} (\lambda - \lambda_i)^{-1} & (\lambda - \lambda_i)^{-2} & \dots & (\lambda - \lambda_i)^{-k} \\ & (\lambda - \lambda_i)^{-1} & & \\ & & \dots & \vdots \\ & & & (\lambda - \lambda_i)^{-1} \end{pmatrix} S_i,$$

The multiplication with the characteristic polynomial allows to cancel the terms with negative exponent.

The resulting expression is a source of eigenvalue – eigenvector relations.

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We express the adjugate with the aid of compound matrices,

$$\text{adj } A \equiv SC_{n-1}(A^T)S.$$

Then we have equality

$$\begin{aligned} P \equiv C_{n-1}(\lambda I - A^T) &= (SVS)G(S\hat{V}^H S) \\ &\equiv (SVS)\chi_A(\lambda)E(S\hat{V}^H S). \end{aligned}$$

The elements of the compound matrix  $P$  are *polynomials* in  $\lambda$  of the form

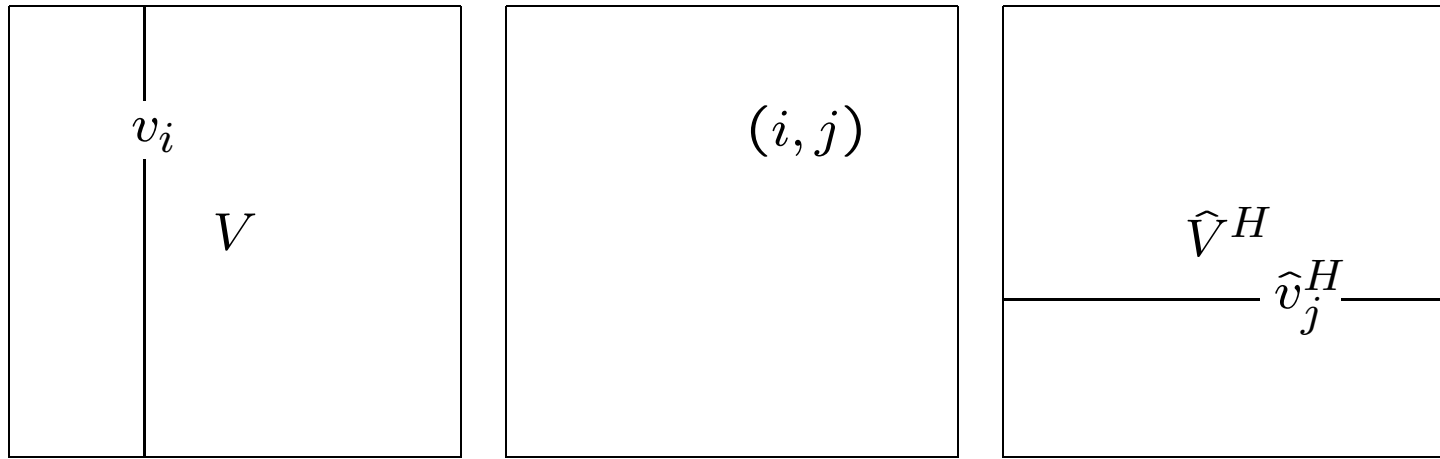
$$p_{ij} = p_{ij}(\lambda; A) \equiv \det L_{ji}, \quad \text{where} \quad L \equiv \lambda I - A.$$

The elements of  $G$  are obviously given by *rational functions* in  $\lambda$ , since

$$G = \chi_A(\lambda) \cdot (\oplus_i E_i).$$

Many terms cancel, the elements of  $G$  are *polynomials*. We divide by the maximal factor and compute the limes  $\lambda \rightarrow \lambda_i$ .

The choice of eigenvectors is based on the non-zero positions  $i, j$  in the matrix (the sign matrices are left out):



Amongst others, the well-known result on eigenvalue – eigenvector relations by Thompson and McEntegert is included. This is one of the basic results used in Paige’s analysis of the finite precision symmetric Lanczos method.

We consider here only the special case of non-derogatory eigenvalues.

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**Theorem:** Let  $A \in \mathbb{K}^{n \times n}$ . Let  $\lambda_l = \lambda_{l+1} = \dots = \lambda_{l+k}$  be a geometrically simple eigenvalue of  $A$ . Let  $k+1$  be the algebraic multiplicity of  $\lambda$ . Let  $\hat{v}_l^H$  and  $v_{l+k}$  be the corresponding left and right eigenvectors with appropriate normalization.

Then

$$v_{jl} \tilde{v}_{i,l+k} = (-1)^{(j+i+k)} \frac{p_{ji}(\lambda_l; A)}{\prod_{\lambda_s \neq \lambda_l} (\lambda_l - \lambda_s)}$$

holds true.

The minus one stems from the sign matrices, the polynomial from the definition of the adjoint as matrix of cofactors and the denominator by division with the maximal factor.

This setting matches every eigenvalue of non-derogatory  $A$ .

Unreduced Hessenberg matrices are non-derogatory matrices. This is easily seen by a simple rank argument. In the following let  $H = H_m$  be unreduced Hessenberg of size  $m \times m$ ,

$$\text{rank}(H - \theta I) \geq m - 1.$$

Many polynomials can be evaluated in case of Hessenberg matrices:

**Theorem:** The polynomial  $p_{ji}$ ,  $i \leq j$  has degree  $(i - 1) + (m - j)$  and can be evaluated as follows:

$$\begin{aligned}
 p_{ji}(\theta; H) &= \begin{vmatrix} \theta I - H_{1:i-1} & & \star \\ & R_{i+1:j-1} & \\ 0 & & \theta I - H_{j+1:m} \end{vmatrix} \\
 &= (-1)^{i+j} \chi_{H_{1:i-1}}(\theta) \prod \text{diag}(H_{i:j}, -1) \chi_{H_{j+1:m}}(\theta).
 \end{aligned}$$

Denote by  $\mathcal{H}(m)$  the set of unreduced Hessenberg matrices of size  $m \times m$ . The general result on eigenvalue – eigenvector relations can be simplified to read:

**Theorem:** Let  $H \in \mathcal{H}(m)$ . Let  $i \leq j$ . Let  $\theta$  be an eigenvalue of  $H$  with multiplicity  $k + 1$ . Let  $s$  be the unique left eigenvector and  $\hat{s}^H$  be the unique right eigenvector to eigenvalue  $\theta$ .

Then

$$(-1)^k \check{s}(i)s(j) = \left[ \frac{\chi_{H_{1:i-1}} \chi_{H_{j+1:m}}(\theta)}{\chi_{H_{1:m}}^{(k+1)}} \right] \prod_{l=i}^{j-1} h_{l+1,l} \quad (1)$$

holds true.

**Remark:** We ignored the implicit scaling in the eigenvectors imposed by the choice of eigenvector-matrices, i.e. by  $\check{S}^T S = I$ .

Among these relations of special interest is the case of index pairs  $(i, m)$ ,  $(1, m)$  and  $(1, m)$ ,  $(1, j)$ :

$$\begin{aligned}
 (-1)^k \check{s}(i) s(m) &= \left[ \frac{\chi_{H_{1:i-1}}(\theta)}{\chi_{H_{1:m}}(\theta)} \right] \prod_{l=i}^{m-1} h_{l+1,l}, \\
 (-1)^k \check{s}(1) s(m) &= \left[ \frac{1}{\chi_{H_{1:m}}(\theta)} \right] \prod_{l=1}^{m-1} h_{l+1,l}, \\
 (-1)^k \check{s}(1) s(j) &= \left[ \frac{\chi_{H_{j+1:m}}(\theta)}{\chi_{H_{1:m}}(\theta)} \right] \prod_{l=1}^{j-1} h_{l+1,l}.
 \end{aligned}$$

These relations are used to derive relations between eigenvalues and *one* eigenvector.

They are also of interest for the understanding of the convergence of Krylov methods, at least in context of Krylov eigensolvers.

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**Theorem:** Let  $H \in \mathcal{H}(m)$ . Let  $\theta$  be an eigenvalue of  $H$ . Then  $\hat{s} = \check{s}$  defined by non-zero  $\check{s}(1)$  and the relations

$$\frac{\check{s}(i)}{\check{s}(1)} = \frac{\chi_{H_{i-1}}(\theta)}{\prod_{l=1}^{i-1} h_{l+1,l}} \quad \forall i \in \underline{m},$$

is (up to scaling) the unique left eigenvector of  $H$  to eigenvalue  $\theta$ .

**Theorem:** Let  $H \in \mathcal{H}(m)$ . Let  $\theta$  be an eigenvalue of  $H$ . Then  $s$  defined by non-zero  $s(m)$  and the relations

$$\frac{s(j)}{s(m)} = \frac{\chi_{H_{j+1:m}}(\theta)}{\prod_{l=j+1}^m h_{l,l-1}} \quad \forall j \in \underline{m},$$

is (up to scaling) the unique right eigenvector of  $H$  to eigenvalue  $\theta$ .

Since the polynomials remain unchanged, merely the eigenvalue moves, this helps to explain convergence behaviour (even in finite precision).



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## Error Analysis Revisited

For simplicity we assume that the perturbed Krylov decomposition

$$M_k = AQ_k - Q_kC_k + F_k$$

is diagonalisable, i.e. that  $A$  and  $C_k$  are diagonalisable.

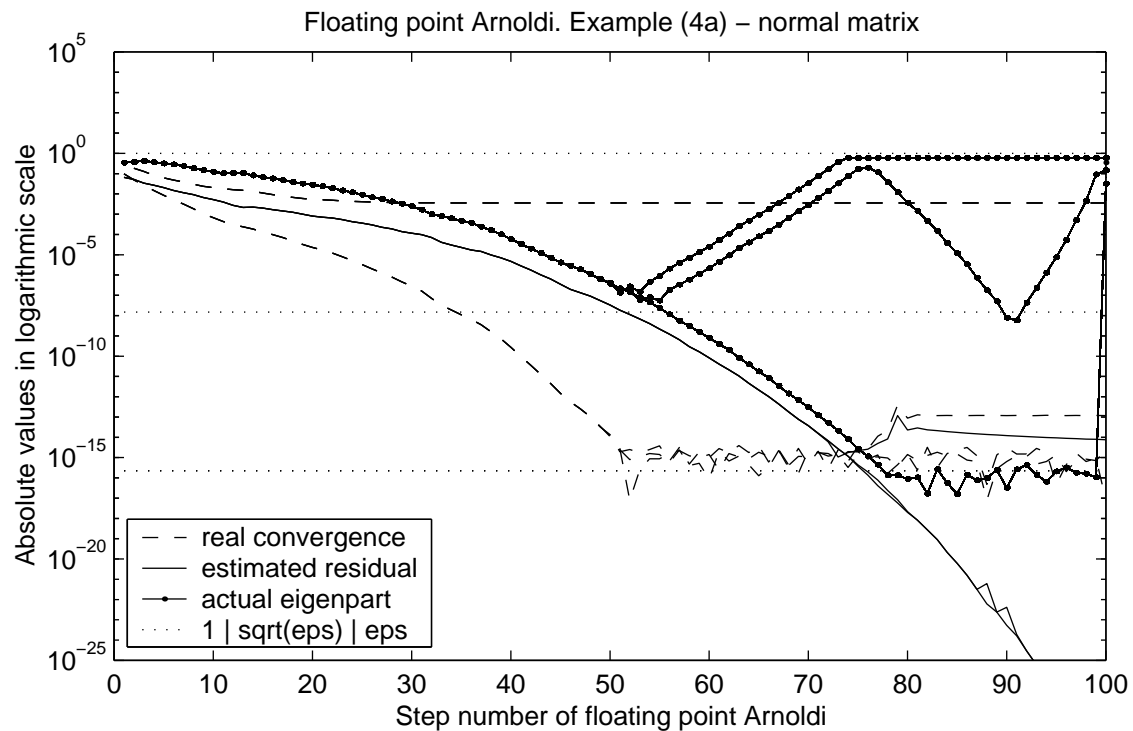
**Theorem:** The recurrence of the basis vectors in eigenparts is given by

$$\widehat{v}_i^H q_{k+1} = \frac{(\lambda_i - \theta_j) \widehat{v}_i^H y_j + \widehat{v}_i^H F_k s_j}{c_{k+1,k} s_{kj}} \quad \forall i, j(, k).$$

This *local error amplification formula* consists of:

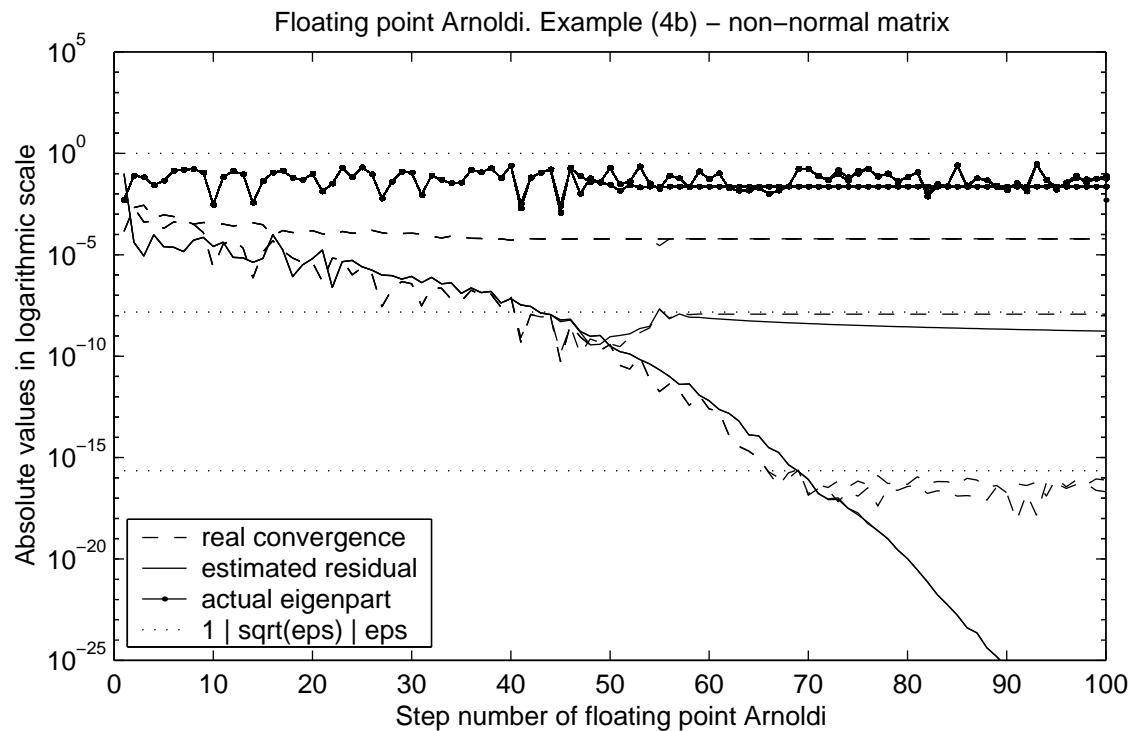
- o the left eigenpart of  $q_{k+1}$ :  $\widehat{v}_i^H q_{k+1}$ ,
- o a measure of convergence:  $(\lambda_i - \theta_j) \widehat{v}_i^H y_j$ ,
- o an error term:  $\widehat{v}_i^H F_k s_j$ ,
- o an amplification factor:  $c_{k+1,k} s_{kj}$ .

$A \in \mathbb{R}^{100 \times 100}$  normal, eigenvalues equidistant in  $[0, 1]$ .



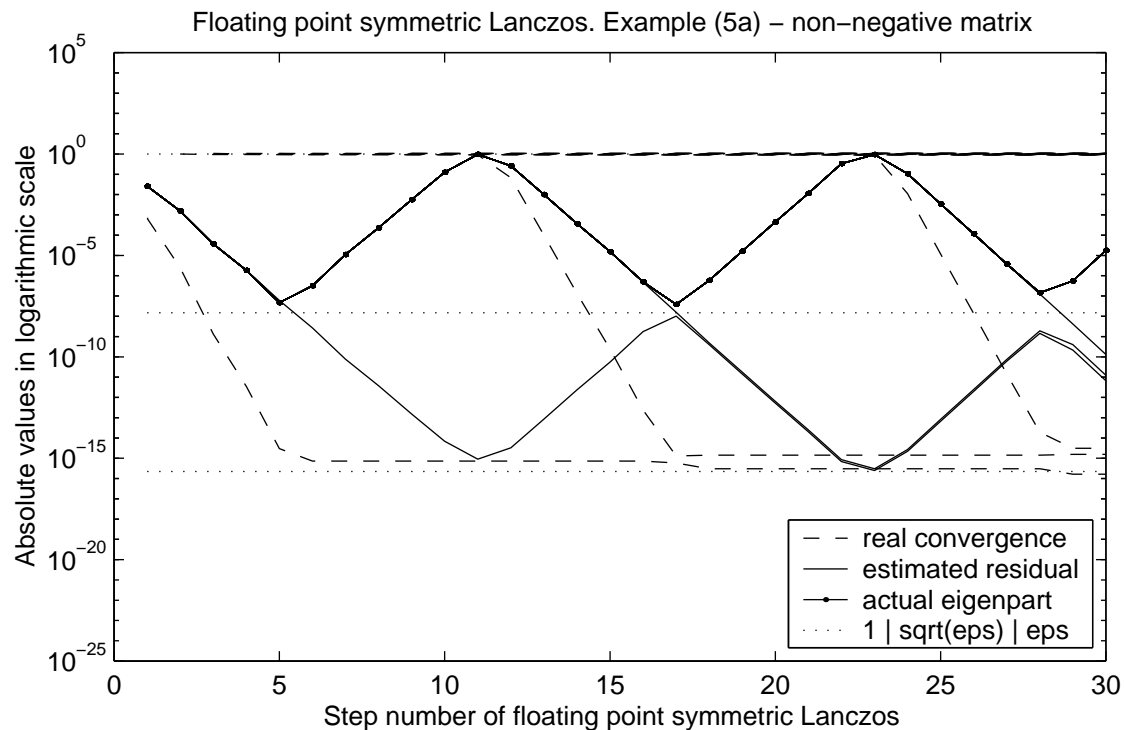
Behaviour of CGS-Arnoldi, MGS-Arnoldi, DO-Arnoldi, convergence to largest eigenvalue.

$A \in \mathbb{R}^{100 \times 100}$  non-normal, eigenvalues equidistant in  $[0, 1]$ .



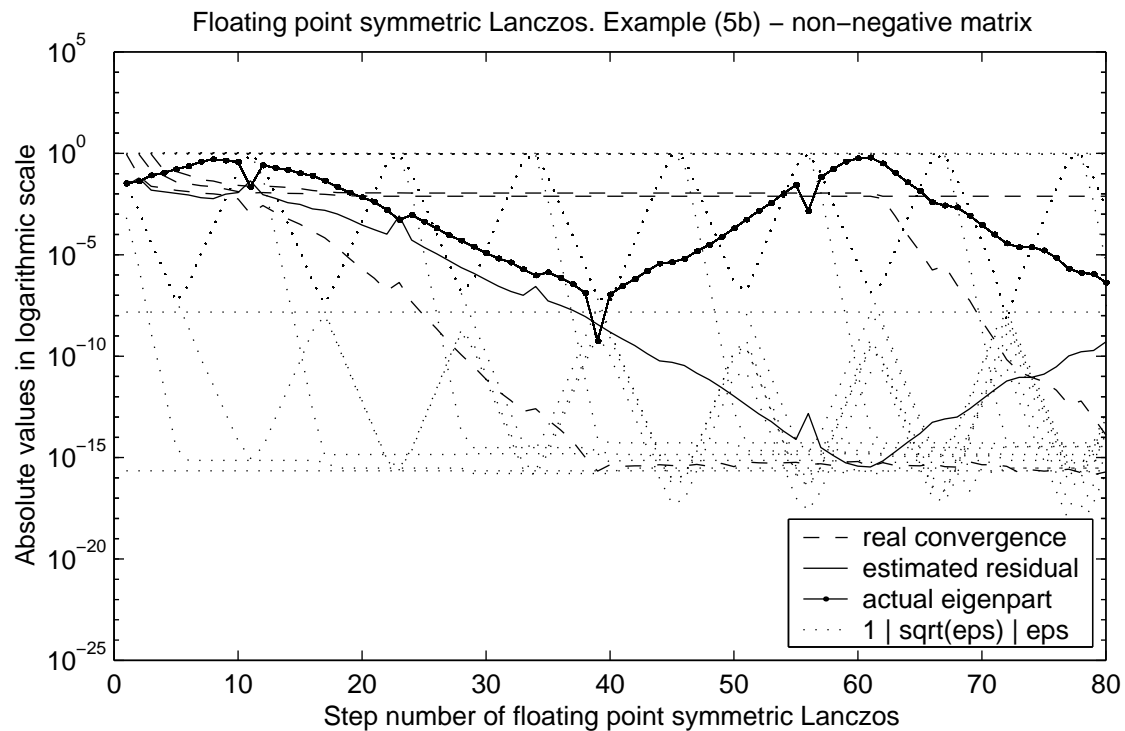
Behaviour of CGS-Arnoldi, MGS-Arnoldi, DO-Arnoldi, convergence to largest eigenvalue.

$A = A^T \in \mathbb{R}^{100 \times 100}$ , random entries in  $[0, 1]$ . Perron root well separated.



Behaviour of symmetric Lanczos, convergence to eigenvalue of largest modulus.

$A = A^T \in \mathbb{R}^{100 \times 100}$ , random entries in  $[0, 1]$ . Perron root well separated.



Behaviour of symmetric Lanczos, convergence to eigenvalue of largest and second largest modulus.

The formula depends on the *Ritz pair* of the actual step. Using the eigenvector basis we can get rid of the *Ritz vector*:

$$I = SS^{-1} = S\check{S}^T \quad \Rightarrow \quad e_l = S\check{S}^T e_l \equiv \sum_{j=1}^k \check{s}_{lj} s_j.$$

**Theorem:** The recurrence between vectors  $q_l$  and  $q_{k+1}$  is given by

$$\left[ \sum_{j=1}^k \frac{c_{k+1,k} s_{kj} \check{s}_{lj}}{\lambda_i - \theta_j} \right] \hat{v}_i^H q_{k+1} = \hat{v}_i^H q_l + \hat{v}_i^H F_k \left[ \sum_{j=1}^k \left( \frac{\check{s}_{lj}}{\lambda_i - \theta_j} \right) s_j \right].$$

For  $l = 1$  we obtain a formula that reveals how the errors affect the recurrence from the beginning:

$$\left[ \sum_{j=1}^k \frac{c_{k+1,k} s_{kj} \check{s}_{1j}}{\lambda_i - \theta_j} \right] \hat{v}_i^H q_{k+1} = \hat{v}_i^H q_1 + \hat{v}_i^H F_k \left[ \sum_{j=1}^k \left( \frac{\check{s}_{1j}}{\lambda_i - \theta_j} \right) s_j \right].$$

**Interpretation:** The size of the deviation depends on the *size* of the *first component* of the *left* eigenvector  $\hat{s}_j$  of  $C_k$  and the *shape and size* of the *right* eigenvector  $s_j$ .

Next step: Application of the eigenvector – eigenvalue relation

$$(-1)^k \check{s}(i) s(j) = \left[ \frac{\chi_{H_{1:i-1}} \chi_{H_{j+1:m}}}{\chi_{H_{1:m}}^{(k+1)}}(\theta) \right] \prod_{l=i}^{j-1} h_{l+1,l}.$$

**Theorem:** The recurrence between basis vectors  $q_1$  and  $q_{k+1}$  can be described by

$$\left[ \sum_{j=1}^k \frac{\prod_{p=1}^k c_{p+1,p}}{\prod_{s \neq j} (\theta_s - \theta_j) (\lambda_i - \theta_j)} \right] \hat{v}_i^H q_{k+1} = \hat{v}_i^H q_1 + \hat{v}_i^H F_k \left[ \sum_{j=1}^k \left( \frac{\check{s}_{1j}}{\lambda_i - \theta_j} \right) s_j \right]$$

A result from polynomial interpolation (Lagrange):

$$\begin{aligned} \sum_{j=1}^k \frac{1}{\prod_{l \neq j} (\theta_j - \theta_l) (\lambda_i - \theta_j)} &= \frac{1}{\chi_{C_k}(\lambda_i)} \sum_{j=1}^k \frac{\prod_{l \neq j} (\lambda_i - \theta_l)}{\prod_{l \neq j} (\theta_j - \theta_l)} \\ &= \frac{1}{\chi_{C_k}(\lambda_i)} \end{aligned}$$

The following theorem holds true:

**Theorem:** The recurrence between basis vectors  $q_1$  and  $q_{k+1}$  can be described by

$$\hat{v}_i^H q_{k+1} = \frac{\chi_{C_k}(\lambda_i)}{\prod_{p=1}^k c_{p+1,p}} \left( \hat{v}_i^H q_1 + \hat{v}_i^H F_k \left[ \sum_{j=1}^k \left( \frac{\check{s}_{1j}}{\lambda_i - \theta_j} \right) s_j \right] \right).$$



Similarly we can get rid of the eigenvectors  $s_j$  in the error term:

$$e_l^T \left[ \sum_{j=1}^k \left( \frac{\tilde{s}_{1j}}{\lambda_i - \theta_j} \right) s_j \right] = \sum_{j=1}^k \left( \frac{\tilde{s}_{1j} s_{lj}}{\lambda_i - \theta_j} \right) = \frac{\prod_{p=1}^l c_{p+1,p} \chi_{C_{l+1:k}}(\lambda_i)}{\chi_{C_k}(\lambda_i)}$$

This results in the following theorem:

**Theorem:** The recurrence between basis vectors  $q_1$  and  $q_{k+1}$  can be described by

$$\begin{aligned} \hat{v}_i^H q_{k+1} &= \frac{\chi_{C_k}(\lambda_i)}{\prod_{p=1}^k c_{p+1,p}} \left( \hat{v}_i^H q_1 + \hat{v}_i^H \sum_{l=1}^k \frac{\prod_{p=1}^l c_{p+1,p} \chi_{C_{l+1:k}}(\lambda_i)}{\chi_{C_k}(\lambda_i)} f_l \right) \\ &= \frac{\chi_{C_k}(\lambda_i)}{\prod_{p=1}^k c_{p+1,p}} \hat{v}_i^H q_1 + \sum_{l=1}^k \left( \frac{\chi_{C_{l+1:k}}(\lambda_i)}{\prod_{p=l+1}^k c_{p+1,p}} \hat{v}_i^H f_l \right). \end{aligned}$$

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Multiplication by the right eigenvectors  $v_i$  and summation gives the familiar result

**Theorem:** The recurrence of the basis vectors of a finite precision Krylov method can be described by

$$q_{k+1} = \frac{\chi_{C_k}(A)}{\prod_{p=1}^k c_{p+1,p}} q_1 + \sum_{l=1}^k \left( \frac{\chi_{C_{l+1:k}}(A)}{\prod_{p=l+1}^k c_{p+1,p}} f_l \right).$$

This result holds true even for non-diagonalisable matrices  $A, C_k$ .

The method can be interpreted as an *additive mixture* of several instances of the same method with several starting vectors.

A *severe deviation* occurs when one of the characteristic polynomials  $\chi_{C_{l+1:k}}(A)$  becomes large compared to  $\chi_{C_k}(A)$ .

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## Open Questions

- o Can Krylov methods be forward or backward stable?
- o If so, which can?
- o Are there any matrices  $A$  for which Krylov methods are stable?
- o Does the stability depend on the starting vector?
- o Are there any **a priori** results on
  - the behaviour to be expected and
  - the rate of convergence?